n‑Type Carbon Nanotubes/Silver Telluride Nanohybrid Buckypaper with a High-Thermoelectric Figure of Merit

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S Supporting Information

[AB](#page-4-0)STRACT: [n-Type therm](#page-4-0)oelectric (TE) materials was made from carbon nanotube (CNT) buckypapers. We used silver telluride $(Ag₂Te)$ to achieve electron injection to the CNTs. The TE characterizations on more than 50 samples show that the CNTs/Ag2Te hybrids exhibit negative Seebeck coefficients (e.g., n-type) from -30 to -228 μ V/K. Meanwhile, the tunneling coupling between the CNTs and Ag₂Te increase the electrical conductance to the range of 10 000–20 000 S/m , which is higher than each single component (CNTs or Ag₂Te). These n-type TE buckypapers are flexible and robust with ZT values of 1−2 orders of magnitude higher than previously reproted for CNT-based TE materials. In addition, the preparation of such buckypapers are very simple compared to a tranditonal inorganic process, without the need for hot pressing or spark sintering. These n-type TE buckypapers can provide important components for fabricating CNT-based flexible TE devices with good conversion efficiency.

KEYWORDS: carbon nanotube, silver telluride, thermoelectric, figure of merit, Seebeck, flexible

NO INTRODUCTION

In the past few decades, thermoelectric (TE) materials have achieved great enhancements in converting waste heat to power.^{1,2} The efficiency of TE materials, figure of merit, ZT, can be expressed as $ZT = S^2 \sigma T / \kappa$, where S is the Seebeck coeffic[ien](#page-5-0)t, σ is the electrical conductivity, T is the absolute temperature, and κ is the thermal conductivity. To obtain advanced efficiency of TE materials, a high Seebeck coefficient, high electrical conductivity, and low thermal conductivity are required. Many efforts on enhancing TE performances in recent years are nanostructuring the materials and, hence, significantly decreasing the thermal conductivity.^{3−6} Thus far, most of the advanced TE materials are reported on inorganic crystals, 7^{-13} which are rigid and brittle. Many ti[mes](#page-5-0), the TE devices made from these inorganic crystals may be subjected to failure u[nder](#page-5-0) continuous thermal cycling because of the mismatch of the thermal expansion among different components. Some strategies have been developed to solve this issue. For example, it has been reported that carbon nanotubes can be used as the interface material to combine with $Bi₂Te₃$ to improve the thermomechanical compliance.¹⁴ Development of flexible TE materials is another attractive approach. Organic materials, such as conducting polymers, w[er](#page-5-0)e reported as flexible TE materials.¹⁵ However, those materials often have low electrical conductivity (<500 S/m), which leads to a low ZT value.^{16,17} Also, it [ma](#page-5-0)y be difficult to expand their application into the high-temperature range, e.g., >500 K.¹⁸

A carbon nanotube (CNT) has many advantages, such as chemi[cal](#page-5-0)ly stable,^{19,20} highly electrically conductive,^{21,22} and

mechanically robust.^{23,24} Owning to these advanced properties as well as abundance of carbon on the Earth, CNT is an important candidate [to b](#page-5-0)uild flexible electronics.25−²⁷ However, it is seldom regarded as a TE material because of its poor Seebeck coefficient $(*60* \mu V/K)^{28}$ and hig[h ther](#page-5-0)mal conductivity of a single tube (\sim 250−10 000 W m⁻¹ K⁻¹ for single tube or aligned arrays).^{29−31} It [has](#page-5-0) been reported that low thermal conductivities can be obtained in CNTs/conductive polymer composites a[nd pr](#page-5-0)ocessed tubes.³² For example, combining CNTs with polyaniline can reduce thermal conductivity to ~0.5 W m⁻¹ K⁻¹.³³ Also, [in](#page-5-0) our previous . reported work, plasma-treated CNT buckypaper shows attractive p-type Seebeck values o[f](#page-5-0) ∼350 μV/K, while the electrical conductivity still maintains an acceptable value (∼1000 S/m).³⁴ Hence, CNT-based materials still hold the promising potential to achieve TE properties comparable to some of the r[ep](#page-5-0)orted traditional inorganic materials, such as BiSbTe^{35,36} or PbTe.^{37–39} On the basis of reported literature, more studies on p-type behavior were reported for TE measur[em](#page-5-0)[en](#page-6-0)t of CN[T-rela](#page-6-0)ted samples.⁴⁰⁻⁴³ Thus far, although there are a few works on n-type CNT-related materials, either a low Seebeck value^{44,45} or a low ele[ctrica](#page-6-0)l conductivity⁴⁶ is obtained and only low ZT values of $\sim 10^{-3}$ are achieved.⁴⁷ Hence, more effort[s sho](#page-6-0)uld be performed to achieve ZT [to](#page-6-0) an acceptable value (e.g., > 0.1), which is comparable to traditio[nal](#page-6-0)

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Figure 1. (A) Optical image, (B) FESEM image, (C and D) TEM images, (E) HRTEM image, and (F) XRD pattern of the CNTs/Ag₂Te hybrid.

inorganic materials. Development of such n-type TE materials is encouraging to provide important components for fabricating CNT-based flexible TE devices.

Herein, we report the preparation of n-type TE materials of silver telluride (Ag_2Te) -decorated single-walled carbon nanotubes (SWCNTs; Sigma-Aldrich; carbon > 90%; ≥77% carbon as SWCNT) (CNTs/Ag₂Te). The sample was prepared via an oil-phase step and then processed into flexible buckypapers. The size of the $Ag₂Te$ particle is around 10 nm. The mass ratio between CNTs and Ag₂Te can be easily tuned by varying the precursor ratio. The TE characterizations on more than 50 samples show that the CNTs/Ag₂Te hybrids exhibit negative Seebeck coefficients (e.g., n-type) from -30 to $-228 \mu V/K$. The electrical conductivity increases for samples with higher amounts of Ag2Te and can reach the range of 10 000−20 000 S/m. The measured thermal conductivity for such hybrid buckypapers can be as low as ~0.7 W m⁻¹ K⁻¹ in the temperature range of 325−525 K. If we do not consider the anisotropic aspect of such buckypaper (e.g., this assumption

allows us to calculate the figure of merit, ZT, based on the measured electric conductivities, Seebeck coefficients, and thermal conductivities), the estimated highest figure of merit, ZT, of CNTs/Ag₂Te hybrids was 0.44 (\pm 5%) at 525 K. Such a ZT value is comparable to some of the traditional TE materials and is higher than that of either pristine CNT or pure $Ag₂Te$.

EXPERIMENTAL SECTION

Synthesis. In a typical synthesis, silver precursor was first prepared by 170 mg (1 mmol) of AgNO₃ dissolving in 10 mL of deionized (DI) water in a flask. Then, 10 mL of toluene with 0.24 mL (1 mmol) of 1 dodecanethiol was added to the solution and well-stirred for 1 h at room temperature for ligand attachment. After that, the stirring was stopped and the solution was self-divided into two phases. Next, the clear aqua phase at the bottom of the flask was removed. The CNTs (1, 3, and 5 mg) were dispersed in toluene via 2 h of ultrasonication. Then, the well-dispersed CNTs were added to the silver precursor solution first. The temperature was increased to 383 K and maintained for 10 min or so. After that, 0.5 mL of TOP-Te solution (1 M) was injected into the as-prepared solution and stirred for 2 min at 383 K.

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The 1 M TOP-Te solution was prepared via dissolving 1.276 g of Te powder into 10 mL of trioctylphosphine. After 2 min of reaction, the flask was quickly quenched in cold water. Next, the sample solution was further dissolved in hexane and then washed by ethanol 2−3 times. The washed solutions were dropped cast onto copper foils and dried in the oven with the temperature of 333 K. After the solvent evaporated, the Ag_2Te/CNT sheets were self-peeled off from the copper foils. To remove the −thiol capping on the samples, the samples were annealed in Ar at 673 K for 1 h. For the Ag_2Te sample, because it was not flexible enough to make a free-standing sheet, the powder was collected and hot-pressed into a pellet in Ar at 673 K.

Characterization Techniques. The morphology and nanostructures of the samples were characterized using X-ray diffraction (XRD; Bruker D8 Powder), field emission scanning electron microscopy (FESEM; JEOL JSM7600), and transmission electron microscopy (TEM; JEOL JEM-2100F). The TEM samples were prepared by dropping the solutions onto copper grids and annealing the whole copper grid in Ar at 673 K. The Hall measurement was carried out using Hall system HL 5500 at room temperature under the van der Pauw mode. The Seebeck coefficient and resistivity were measured from 323 to 523 K using the ZEM-3 Seebeck meter under a helium environment (see Figure S1 of the Supporting Information). The samples were cut into a rectangular piece with a dimension of 10×4 mm² and fixed on a glass substrate by applying silver paste to both ends. The samples [for](#page-4-0) [thermal](#page-4-0) [conductivity](#page-4-0) [measurements](#page-4-0) [we](#page-4-0)re cut into $12.7 \times 12.7 \text{ mm}^2$ (CNTs/Ag₂Te hybrid and pure CNT buckypaper) or polished into the round pellet with the diameter of 12.7 mm $(Ag_2Te$ pellet) and were measured by laser flash (NETZSCH MicroFlash LFA457).

■ RESULTS AND DISCUSSION

Five types of samples, including pure $Ag₂Te$, pure CNTs, and three $CNTs/Ag_2Te$ hybrids, were prepared for characterization and TE property measurements. The $CNTs/Ag_2Te$ hybrid samples were prepared by injecting TOP-Te into oil-phase solution containing Ag precursor, CNT, and surfactant 1 dodecanethiol. The weight ratios of CNT Ag_2Te can be adjusted by changing the amount of CNT and estimated on the basis of the amount of precursor added. The estimated weight ratios of CNT/Ag_2Te of three hybrid composites were 1:17, 3:17, and 5:17 and named as CA-1, CA-2, and CA-3, respectively. Figure 1A is the optical image of the buckypaper of sample CA-1. It was prepared via drop casting the washed CNTs/Ag2Te hybri[d](#page-1-0) solution onto copper foil and peeling off the free-standing paper after drying. The buckypapers were then annealed at 673 K. The obtained buckypaper was around 50 μ m in thickness and mechanically robust upon repeated bending. The low-magnification FESEM image of hybrid sample CA-1 (Figure 1B) indicates that the CNT bundles were entangled randomly. Panels C−E of Figure 1 are transmission electron [m](#page-1-0)icroscopy (TEM) images that show the morphology and structure of the hybrid sample C[A-](#page-1-0)1. It can be found that most $Ag₂Te$ particles attached on the CNT walls are around 10 nm. Small particles may agglomerate and form large particles during the annealing process. One large particle of around 25 nm, which is denoted by the arrow, can also be observed (Figure 1C). Additionally, we can find that the diameter of a single CNT is <1 nm, but the CNTs may selfassemble into bundles [\(F](#page-1-0)igure 1D) because of the van der Waals force. The diameters of CNT bundles shown in panels C and D of Figure 1 are around [1](#page-1-0)0 nm. Figure 1E is a highresolution TEM (HRTEM) image of an Ag_2Te particle on CNT walls. It i[nd](#page-1-0)icates that the Ag_2Te part[icl](#page-1-0)e is singlecrystalline. The representative XRD pattern of CA-1 is shown in Figure 1F, which matches the monoclinic Ag_2Te phase [Joint

Committee on Powder Diffraction Standards (JCPDS) number 34-0142]. There is no detectable impurity phase.

For comparison purposes, pure Ag₂Te and pure CNT samples were also prepared. Ag_2Te was synthesized by a similar method as that of CNTs/Ag2Te hybrids but without adding CNTs. Figure 2A shows the XRD pattern of pure Ag_2Te . Peaks

Figure 2. (A) XRD pattern of pure Ag_2Te . (B) TEM image of pure $Ag₂Te$ nanoparticles. (C) SEM image of pure CNT random networks.

corresponding to the monoclinic Ag2Te phase (JCPDS number 34-0142) can be observed. There is a reversible phase transition of Ag₂Te at 418 K from the β phase with a monoclinic structure to the α phase with a face-centered cubic (fcc) structure. Although the samples were annealed at 673 K after synthesis, no impurity phase, except for the monoclinic Ag₂Te phase, is detectable according to the XRD result. Thus, we believe that Ag₂Te can be easily converted back from the fcc phase to the monoclinic phase when the temperature returns to lower temperatures (<418 K). Figure 2B shows the TEM image of pure Ag2Te particles. The nanoparticles are uniform with the size of ∼5 nm. Figure 2C is the FESEM image of pure CNT buckypaper. It shows a random network of CNT bundles. The average diameter of CNT bundles is around 10 nm.

For TE property measurements, $CNTs/Ag_2Te$ hybrid buckypapers with a thichness of ∼50 μm were cut into proper size for different measurements. Pure CNTs were also prepared into flexible buckypaper for TE measurements. Pure Ag_2Te was rigid and brittle; therefore, it cannot be prepared into a freestanding sheet. Hence, the Ag_2Te pellet was prepared via hot press at 673 K.

Electrical conductivity values of three CNTs/Ag₂Te hybrid buckypapers (CA-1, CA-2, and CA-3) are shown in Figure 3A. It shows that samples with a higher Ag_2Te content exhibit higher electrical conductivities at the same temperature ra[ng](#page-3-0)e. CA-1 shows the highest electrical conductivity of 23 500 S/m at 381 K. The shape of the electrical conductivity curves are similar for these three samples within the measurement temperature of 325−550 K. The electrical conductivity of

Figure 3. (A) Electrical conductivity, (B) Seebeck coefficient, (C) power factor, and (D) thermal conductivity of CNTs/Ag₂Te hybrid buckypapers, pure Ag2Te pellet, and pure CNT buckypaper.

CNTs/Ag2Te hybrids drops suddenly at a temperature close to the phase transition point of Ag_2Te (i.e., 418 K). Below the phase transition point, the electrical conductivity becomes higher at an increased temperature, which shows a typical semiconducting behavior. At a temperature range of $T > 418$ K, the electrical conductivity decreases with the increase of the temperature, which becomes a metallic/heavily doped semiconducting behavior. A pure Ag₂Te pellet and pure CNT buckypaper were also measured for comparison. The electrical conductivity curve of the pure Ag2Te pellet displays a similar trend to that of CNTs/Ag₂Te hybrids. It increases at an increased temperature below the phase transition point and decreases at an increased temperature above the phase transition point. The electrical conductivity of CNT buckypaper is 3500−5000 S/m in the measured temperature range. It decreases at increased temperatures, which shows a metallic/ heavily doped semiconducting behavior. The electrical conductivity of CA-1 is higher than both pure CNT buckypaper and pure $Ag₂Te$ pellet, which shows the synergism in the CNTs/Ag2Te hybrid structure similar to the previous reported effect for binary nanocrystal superlattices.⁴⁸ The higher electrical conductivity of CA-1 compared to pure CNT buckypaper is possibly related to the dop[ing](#page-6-0) effect. The electrical properties of all CNTs/Ag₂Te hybrid samples and the pure Ag₂Te pellet can be further studied with Hall measurements at room temperature (Table 1). The mobility, μ , can be calculated via $\sigma = n\mu e$, where σ is the electrical conductivity, *n* is the carrier concentration, and e is the elementary charge. For the three CNTs/Ag₂Te hybrids, the carrier concentration is

Table 1. Electrical Conductivity, Carrier Concentration, and Mobility of $CNTs/Ag_2Te$ Hybrids and Pure Ag_2Te

	σ (S/m)	$n \text{ (cm}^{-3})$	μ (cm ² V ⁻¹ s ⁻¹)
$CA-1$	1.91×10^{4}	6.59×10^{18}	181.1
$CA-2$	8.10×10^{3}	2.05×10^{18}	246.7
$CA-3$	2.11×10^3	1.09×10^{18}	120.7
Ag ₂ Te	1.85×10^{4}	9.10×10^{19}	12.7

higher for samples with more Ag_2Te . The sample CA-1, which contains most Ag2Te, has the largest carrier concentration of 6.591 \times 10¹⁸ cm⁻³ among the three CNTs/Ag₂Te hybrids. This observation indicates that, during the Ag_2Te decoration process, electrons are injected from Ag_2Te to CNTs. With the increase of the Ag_2Te content, more electrons exist in the system; hence, the carrier concentration increases. The mobilities of CNTs/Ag2Te hybrids show a non-monotonic trend. This is because of the fact that, in these hybrid samples, $Ag₂Te$ plays multiple roles, including network element, electron donor, impurity, etc. Therefore, their impact on the electronic properties is complex. For example, as network elements, they provide extra transport channels, helping to enhance electron mobility, while those decorated onto CNTs scatter electrons. Thus, a combination of these interactions between composites and carriers possibly gives rise to a non-monotonic mobility trend.49,50 Additionally, we found that the mobilities of CNTs/ Ag₂Te hybrids (120–250 cm² V⁻¹ s⁻¹) are higher than pure Ag₂T[e \(12](#page-6-0).702 cm² V⁻¹ s⁻¹) by 1 order of magnitude. Although the electron scattering caused by the $CNTs/Ag_2Te$ interfaces may be stronger than the Ag_2Te/Ag_2Te interfaces, the change of the electron mean free path (MFP) should also be considered. Because Ag_2Te is a nanoparticle, the electron MFP in $Ag₂Te$ is shorter than the size of the particle (i.e., 10 nm) and electron transport is limited by the size of the nanoparticle. However, MFP can be much longer in CNTs, considering the length of CNTs (hundreds of nanometers to several micrometers) and their exotic transport property along the tube. Thus, in hybrid samples, the average MFP is enhanced with CNTs as part of the network and mobilities are larger than that of pure $Ag₂Te$.

We mention that the effect of the irregularities in the nanoparticle positions and in the strengths of the tunneling coupling is different for metallic and insulating samples. If the coupling between the particles is sufficiently strong and the system is well-conducting, the irregularities are not very important.⁵¹ In contrast, irregularities become crucial in the limit of low coupling, where the system is an insulator. The key

parameter that determines most of the physical properties of inhomogeneous samples is the dimensional tunneling conductance g between neighboring particles. Samples with $g > 1$ exhibit metallic transport properties, while those with $g < 1$ show insulating behavior. All of our samples belong to the insulating side of the metal−insulator transition. Therefore, the charge on each particle becomes quantized as in the standard Coulomb blockade behavior. In this case, an electron has to overcome an electrostatic barrier of the order E_c to hop onto a neighboring particle, with E_c being the Coulomb energy of a single particle. At applied bias voltages and temperatures large enough to overcome local Coulomb energy costs, transport occurs by sequential tunneling between neighboring dots.⁵² At small bias and low temperatures, sequential tunneling is suppressed by the Coulomb blockade. In this re[gim](#page-6-0)e, conduction involves higher order tunneling processes, the socalled co-tunneling events, which can transport a charge over distances of several particles without incurring the full Coulomb energy costs. The above physics is valid for low temperatures, below 380 K. At higher temperatures, above 380 K, our samples show metallic behavior, meaning that the temperature reaches the size of the Coulomb gap. In this case, the number of electrons in the conducting band increases exponentially with the temperature, leading to metallic behavior: the higher the temperature, the lower the conductivity. This behavior is the consequence of the fact that the number of phonons increases with the temperature, leading to more difficult electron propagation through the sample because of many scattering events.

Seebeck coefficients of three CNTs/Ag₂Te hybrid buckypaper samples display n-type properties (negative Seebeck coefficient) (Figure 3B). Generally, the absolute values of Seebeck coefficients of three CNTs/Ag₂Te hybrid samples are larger at higher te[mp](#page-3-0)eratures in the temperature range of ∼418−550 K. Among them, CA-1 exhibits the highest Seebeck value of 228 μ V/K at 537 K. The two other hybrid samples, CA-2 and CA-3, also achieve Seebeck values of >100 $\mu\bar{V}/K$ when the temperature is higher than 480 K. The maximum Seebeck value obtained for the Ag₂Te pellet is 157 μ V/K at 421 K, which is ∼100 K lower than that of CA-1. The CNT buckypaper shows a lower p-type value (\sim 30−50 μ V/K) within the measured temperature range. The maximum Seebeck value of CA-1 is ~45% higher than that of the pure Ag₂Te pellet and 4 times higher than that of pure CNT buckypaper. The enhancement of the Seebeck value for CA-1 may result from the change of the density of states because of the fact that the Ag2Te dopants may generate some impurity levels in the band structures.

In combination of the Seebeck coefficient and electrical conductivity, the power factors of the $CNTs/Ag_2Te$ hybrids, pure Ag_2Te , and pure CNTs are calculated (Figure 3C). Because of the largest Seebeck values and highest electrical conductivity, sample CA-1 presents the largest power f[ac](#page-3-0)tor compared to those of all of the other samples. The highest power factor of CA-1 is 579 μ W m⁻¹ K⁻² at 537 K, which is 2 times that of the pure Ag₂Te pellet value and 80 times that of the pure CNT buckypaper value in the same temperature range.

Thermal conductivities (Figure 3D) of selected samples, CA-1, pure Ag_2Te , and pure CNTs were measured by laser flash. Sample CA-1 exhibits a low ther[m](#page-3-0)al conductivity of ∼0.7 W $m^{-1}K^{-1}$, which is lower than that of the pure Ag₂Te pellet and pure CNTs buckpaper in the same temperature range. Thermal conductivity, κ , can be expressed as $\kappa = \kappa_e + \kappa_L$, where κ_e and κ_L are electron thermal conductivity and lattice thermal conductivity, respectively. In the absence of inelastic processes, the electron thermal conductivity can be estimated using Wiedemann–Franz relation $\kappa_e = L\sigma T$. Because the Lorenz number cannot be accurately determined, the accurate value of κ_e cannot be confirmed. However, if we assume $L = 2.44 \times 10^{-8}$ W Ω K⁻² and take the point of 525 K as an example, the electron thermal conductivity of CA-1, $Ag₂Te$, and CNT can be estimated as 0.14, 0.1, and 0.05 W $m^{-1} K^{-1}$, respectively. The lattice thermal conductivity can be calculated then as 0.55, 0.68, and 1.03 W m[−]¹ K[−]¹ , respectively. Although sample CA-1 has the highest electron thermal conductivity because of the largest electrical conductivity, it exhibits the lowest lattice thermal conductivity mainly attributed to the strong phonon scattering of the hybrid structure. It is worth noting here that the $Ag₂Te$ sample is a dense pellet prepared by hot press at 673 K, while the CA-1 and pure CNTs samples are loosely packed buckypapers. The structural difference will also affect the lattice thermal conductivity. In general, the total thermal conductivity shows the non-monotonic temperature behavior because of competition between two different mechanisms. At high temperatures, the number of phonons is large leading to larger thermal conductivity.

If we do not consider the anisotropic aspect of such buckypaper, the figure of merit, ZT, can be calculated on the basis of the measured electric conductivities, Seebeck coefficients, and thermal conductivities. The estimated highest figure of merit, ZT, is 0.44 at 525 K on the sample CA-1. It is more than 2 times higher than that of pure $Ag₂Te$ and 100 times higher than that of CNT. Although the TE performance of $Ag₂Te-based materials may be further enhanced with more$ efforts (e.g., tuning particle size⁵³), they are rigid and may have different working temperatures.⁵⁴ Hence, the CNTs/Ag₂Te hybrid holds niche applications [be](#page-6-0)cause it is flexible and robust.

■ CONCLUSION

In summary, we successfully synthesized n-type $CNTs/Ag_2Te$ hybrids via a solvothermal approach. The Ag_2Te particles attached on CNT bundles have a small size of ∼10 nm, and these $CNTs/Ag₂Te$ hybrids can be prepared into flexible buckypaper. The hybrid sample CA-1 shows a higher electrical conductivity, higher Seebeck coefficient, and lower thermal conductivity than those of pure Ag₂Te and pure CNTs in the temperature range of 325−525 K. These enhancements may be attributed to the doping effect, more impurity levels, and stronger phonon scattering, respectively. Hence, its maximum ZT value of 0.44 is 2 times higher than that of the pure Ag_2Te pellet and 100 times higher than that of pure CNT buckypaper. This result shows that $CNTs/Ag_2Te$ hybrids can be a promising new type n-type TE material with mechanical flexibility and acceptable ZT values.

■ ASSOCIATED CONTENT

6 Supporting Information

Schematic of the ZEM-3 setup (Figure S1) and thermal diffusivity and specific heat capacity results (Figure S2). This material is available free of charge via the Internet at http:// pubs.acs.org.

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Notes

The authors declare no competing financial interest.

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